Abstract

In the space in an atom, when a moving electron gains velocity, its kinetic energy increases but its total energy decreases.

From this fact, in the space in an atom, we need to revise the relational expression of special relativity: \( E^2 = c^2 p^2 + E_0^2 \)

In this paper we will obtain a new relational expression which takes the place of that of special relativity. Further, when we calculate the radius of a hydrogen atom with the new relational expression, we can obtain not only Bohr radius but the size of an atomic nucleus (i.e. proton).

Moreover, we will explain without the help of quantum mechanics, why an electron moving around an atomic nucleus is not absorbed in the nucleus.

I. Introduction

One of the important equations of special relativity is:
\[ E^2 = c^2 p^2 + E_0^2 \]  (I. 1)

\( E \) means total energy of an object or a particle. \( E_0 \) means rest mass energy.

When particles moving in the macro space gain velocity, their kinetic energy and total energy increase. From this, we can obtain \( E > E_0 \).

In regard to an electron in an atom, however, when it gains velocity its kinetic energy increases but its total energy decreases. From this it is predictable that in an atom, \( E < E_0 \).

In the next chapter we consider whether Eq. (I . 1), which we can obtain in the macroscopic space, can be obtained in an atom.
### II. Relationships between energy and momentum of an electron in a hydrogen atom

First of all, let us review the energy of an electron in a hydrogen atom.

Suppose the atomic nucleus is at rest because it is heavy. Consider the situation where electron (electric charge $-e$, mass $m$) is circulating at speed $v$ along the orbit (radius $r$) whose center is the nucleus.

We can obtain an equation as follows:

$$\frac{mv^2}{r} = e^2/(4\pi\varepsilon_0 r)$$  \hspace{1cm} (II. 1)

Therefore:

$$\frac{mv^2}{2} = (1/2)e^2/(4\pi\varepsilon_0 r)$$  \hspace{1cm} (II. 2)

Meanwhile, the potential energy $V(r)$ of electron is:

$$V(r) = -e^2/(4\pi\varepsilon_0 r)$$  \hspace{1cm} (II. 3)

Since the right side of Eq.(II. 2) is $-1/2$ times of the potential energy, we have:

$$2\left(\frac{mv^2}{2}\right) = -V(r)$$  \hspace{1cm} (II. 4)

Therefore, relationships between the total energy $E$, kinetic energy $K$ and potential energy are:

$$E = V(r) + K$$

$$= -2K + K$$

$$= -K$$

$$= V(r)/2$$  \hspace{1cm} (II. 5)

It is known that, as to an electron, the absolute value of the total energy and that of the kinetic energy are equal to each other.

The total energy of an electron is equal to the half of the potential energy.

On the basis of the equation above, we will obtain relationships between energy and momentum of an electron in a hydrogen atom referring to a textbook on special relativity.  \(^{1}\)

In classical mechanics, the increment of kinetic energy corresponds to the work done by external forces, and we have

$$dK = Fdx = (dp/dt)dx = vdp$$  \hspace{1cm} (II. 6)

In regard to particles moving in the macroscopic space, the increment of total energy and that of kinetic energy is equal if potential energy doesn’t change, i.e.,

$$dE = dK$$  \hspace{1cm} (II. 7)

Where

$$dE = vdp$$  \hspace{1cm} (II. 8)

In regard to an electron moving in an atom, however, the decrement of total energy is equal to the increment of kinetic energy, as is clear from Eq. (II. 5):

$$-dE = dK$$  \hspace{1cm} (II. 9)

From this and Eq. (II. 6), we have

$$dE = -vdp$$  \hspace{1cm} (II. 10)
We can obtain Eq. (I. 1) from Eq. (II. 8). However, the relationships between the energy and momentum of an electron in a hydrogen atom have to be obtained from Eq. (II. 10).

In classical mechanics,
\[ m = \frac{p}{v} \quad \text{(II. 11)} \]

Also, from special relativity,
\[ m = \frac{E}{c^2} \quad \text{(II. 12)} \]

From Eq.(II. 11) and Eq.(II. 12):
\[ E = c^2 \frac{p}{v} \quad \text{(II. 13)} \]

Next, we multiply together the left and right sides of the two equations (II. 10) and (II. 13):
\[ EdE = - c^2 p dp \quad \text{(II. 14)} \]

Then, we integrate this:
\[ E^2 = - c^2 p^2 + E_0^2 \quad \text{(II. 15)} \]

where \( E_0^2 \) is a constant of integration, written explicitly as the square of some constant energy.
III. The meaning of energy $E$

What does energy $E$ mean in Eq.(II.15)?

In classical mechanics we are never concerned with absolute energies but only with energy differences.

In this paper, however, we consider absolute quantity of an electron’s energy.

In existing theories, it is when we separate electrons from an atomic nucleus infinitely and put them at rest that we regard total energy of electrons as zero. Total energy in Eq. (II.5) is the value obtained from that point of view.

However, even if we put electrons at rest in an infinitely distant place, the absolute energy of electrons is essentially not zero. Electrons should have rest mass energy $E_0$.

Considering this and Eq. (II.5), we define total energy of electrons in an atom $E_{ab}$ in absolute sense as below:

$$E_{ab} = E_0 + V(r) + K$$

$$= E_0 - 2K + K$$

$$= E_0 - K$$

$$= E_0 + V(r)/2$$

(III.1)

We regard $E$ in Eq. (II.15) as total energy $E_{ab}$ in Eq. (III.1), not as total energy in (II.5).

Then, Eq. (II.15) is expressed as:

$$E_{ab}^2 + c^2 p^2 = E_0^2$$

(III.2)

This equation shows the relationships between energy and momentum of an electron in a hydrogen atom. We have got a different result from special relativity’s formula (I.1).
IV. Orbital radius and energy of an electron in a hydrogen atom

In this chapter we will consider whether we will have a new advance in physics from Eq. (III. 2) that we obtained in the previous chapter.

According to classical quantum theory, classical electron radius $a_n$ and energy $E_n$ of a hydrogen atom is given in the following:

$$a_n = \frac{4\pi\varepsilon_0 (\hbar/2\pi)^2}{(m_0e^2)} n^2$$  \hspace{1cm} (IV. 1)

$$E_n = -\frac{e^2}{(8\pi\varepsilon_0 a_n)} \quad (n = 1, 2, \ldots)$$  \hspace{1cm} (IV. 2)

Here, $n$ is a principal quantum number.

From the view point of this paper, what will we get as the value of $a_n$ and $E_n$?

First, from total energy of an electron defined in Eq. (III. 1) and from Eq. (III. 2);

$$\left( E_0 - K \right)^2 + c^2 p^2 = E_0^2$$  \hspace{1cm} (IV. 3)

Therefore

$$p = (2E_0K - K) / c$$  \hspace{1cm} (IV. 4)

By the way, Bohr’s quantum condition is;

$$p \cdot 2\pi n = 2\pi n(h/2\pi)$$  \hspace{1cm} (IV. 5)

Substituting the value in Eq. (IV. 4) for the momentum of Eq. (IV. 5);

$$\left( 2E_0K - K \right) / c = n(h/2\pi)$$  \hspace{1cm} (IV. 6)

Squaring both sides and substituting the value of the right side of Eq. (II. 2) for the kinetic energy;

$$\left[ 2m_0c^2 (1/2c)^2 / (4\pi\varepsilon_0 r_n) - (1/2c^2 e^2 / (4\pi\varepsilon_0 r_n)^2) \right] r_n^2 = n \cdot c^2 (h/2\pi)^2$$  \hspace{1cm} (IV. 7)

Solving this referring to $r_n$;

$$r_n = \left( 1/4 \right) e^2 / (4\pi\varepsilon_0 m_0c^2) + (4\pi\varepsilon_0 (h/2\pi)^2 / (m_0 e^2)) n^2$$

$$= (r_e/4) + (1/r_e)(\lambda_e/2\pi)^2 n^2$$

$$= (r_e/4) + n a_B$$  \hspace{1cm} (IV. 8)

Here, $r_e$ is classical electron radius, and $\lambda_e$ is compton wave length of electrons.

They are given in the following;

$$r_e = e^2 / (4\pi\varepsilon_0 m_0c^2)$$  \hspace{1cm} (IV. 9)

$$\lambda_e = h / (m_0c)$$  \hspace{1cm} (IV. 10)

To the radius $r_n$ obtained in this paper[(IV. 8)], the term $r_e/4$ is added besides the value obtained from classical quantum theory (IV. 1).

Further, when $n=1$, the radius is:

$$r_1 = (r_e/4) + a_B$$  \hspace{1cm} (IV. 11)

(Note that $a_B$ is Bohr radius)

In addition, substituting $r_n$ in Eq. (IV. 8) for Eq. (IV. 2):

$$E_n = -\frac{e^2}{(8\pi\varepsilon_0 r_n)} \quad (n = 1, 2, \ldots)$$  \hspace{1cm} (IV. 12)
V. Conclusion

1. In the macro space we can obtain Eq.(I. 1) of special relativity. In the space of a hydrogen atom, however, we can obtain Eq.(III . 2).

2. It is natural to think that the term $r_e/4$ added newly in Eq. (IV. 8) has something to do with the radius of an atomic nucleus (i.e. proton).

   In addition, from Eq. (IV. 11), we have found atomic orbital radius that Bohr’s classical quantum theory predicts is not the distance from the center of an atomic nucleus to the orbit, but the distance from the surface of an atomic nucleus to the orbit.

   In Bohr radius the size of the atomic nucleus is not taken into consideration.

3. When an electron with the rest mass energy $E_0$ outside an atom is taken in an atom, the total energy of an electron decreases. Suppose the decreased energy is $-E'$. In this case, $E'/2$ is transformed into the kinetic energy of an electron, and the rest $E'/2$ is emitted as a photon outside an atom.

   That is:
   $$-E' + K + (\hbar/2\pi) \omega = 0$$  \hspace{1cm} (V. 1)

   (Note that $(\hbar/2\pi)\omega$ is a photon’s energy)

   Further, considering Eq. (III . 1), $-E'$corresponds to the potential energy of an electron.

   That is:
   $$-E' = V(r)$$  \hspace{1cm} (V. 2)

   Therefore
   $$-E' + K + (\hbar/2\pi) \omega = V(r) + K + (\hbar/2\pi) \omega = 0$$  \hspace{1cm} (V. 3)

   From this, we can predict the existence of the lower limit of an electron’s total energy $E_{ab}$ even in the classical mechanics.

   From Eq. (III . 1), the state where the potential energy has consumed all the rest mass energy is:
   $$E_0 - 2K = 0$$  \hspace{1cm} (V. 4 )

   In this case the kinetic energy of an electron becomes $m_e c^2 / 2$. Therefore, the relations between $E$ (the total energy of an electron in existing theories) and $E_{ab}$ is:
   $$E = -m_e c^2 / 2 = -E_{ab}$$  \hspace{1cm} (V. 5)

   This is the minimum value of an electron’s total energy considered in classical mechanics.

   Substituting this for $E_n$ in Eq. (IV. 12), we can obtain the following value $r$ as the distance of closest approach:
   $$r = e^2 / (4\pi\epsilon_0 m_e c^2) = r_e$$  \hspace{1cm} (V. 6)

   From Eq. (IV. 8), the radius of an atomic nucleus can be considered $r_e/4$. Thus, in the prediction based on classical mechanics, it is clear that an electron is not absorbed in an atomic nucleus.
Furthermore, the distance $r_e$ agrees with the distance of closest approach of $\alpha$-particle in Rutherford scattering\(^2\).

**Acknowledgments**

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**References**


2) E.RUTHERFORD, Phil. Mag. 37 537, 1919